ABSTRACT: The first vapor liquid equilibrium (VLE) measurements for the binary systems of ammonia (NH₃) and three imidazolium-based ionic liquids (ILs) have been successfully measured using a gravimetric microbalance. ILs 1-butyl-3-methylimidazolium hexafluorophosphate ([C₄C₁im][PF₆]), 1-butyl-3-methylimidazolium tetrafluoroborate ([C₄C₁im][BF₄]), and 1-ethyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide ([C₂C₁im][NTf₂]) were measured at temperatures of 283.15, 298.15, 323.15, and 348.15 K and at pressures up to 0.7 MPa using the new Hiden XEMIS gravimetric microbalance. The VLE data were correlated using the Peng–Robinson equation of state and the Non-Random Two Liquid (NRTL) activity coefficient models. Both models are in excellent agreement with the experimental data. The Fickian diffusivities of NH₃ in imidazolium-based ILs were obtained fitting experimental concentration to the one-dimensional (1D) mass diffusion equation, and found to be about 3 to 5 times lower than the diffusion of NH₃ in water (H₂O). The semitheoretical Stokes–Einstein equation was used to model diffusivities and to obtain the diffusion radius of NH₃ in imidazolium-based ILs.

1. INTRODUCTION

Ionic liquids (ILs) are liquid salts composed entirely of ions and defined to have a melting point below 373 K. ILs differ from traditional salts because the charge delocalization, the size of the ion, and/or the ion asymmetry prevent the formation of stable crystals. This unique behavior for ILs categorizes them as a new class of material with many unique properties such as low vapor pressure, good electrochemical and thermal stability, and high solvation capability. Novel characteristics of ILs have led to extensive research to explore the feasibility of ILs in various applications such as gas separation, absorption–refrigeration, and catalysis which requires fundamental thermodynamic knowledge (e.g., vapor liquid equilibria) for ILs with other substances.

Ammonia (NH₃) is one of the most important and largest volume industrial chemicals produced today. Most of the research involving NH₃ and ILs has been directed at reducing NH₃ emissions or to improve the NH₃–H₂O absorption–refrigeration cycle. In 2007, Yokozeki and Shiflett pioneered the investigation of NH₃ and IL mixtures by measuring the NH₃ solubility in eight imidazolium-based ILs suggesting that NH₃ + IL might be an alternative refrigerant–absorbent pair for the NH₃–H₂O absorption refrigeration cycle. Huang et al. showed the solubility of NH₃ in guanidinium-based ILs was comparable with that of imidazolium-based ILs and suggested the interactions between NH₃ and the IL were dominated by the cation. Shi and Maginn using Monte Carlo simulations showed that the basic nitrogen of NH₃ associates with the acidic hydrogen that is attached to the C(2) carbon of the imidazolium ring and the anion has little effect on the solubility of NH₃ for [C₄C₁im][NTf₂]. Li et al. also found the length of the alkyl chain has an impact on NH₃ sorption in imidazolium-based ILs. For example, the NH₃ solubility in imidazolium ILs increased with longer alkyl chains on the cation ([CₙC₁im][BF₄], n = 2, 4, 6, 8). Tomida et al. recently measured the
solubility of NH₃ in [C₄C₁im][PF₆] over a wide temperature range (283.15 to 363.15 K) and at pressures up to 3 MPa.⁸

The NH₃ absorption capacity was also measured in various hydroxyl functionalized ILs: [N₁₁₁₂C₁OH][NTf₂],⁹⁻¹¹ [(HOCC₂)C₃im][BF₄],⁹⁻¹² [(HOCC₂)C₅im][DCA],¹⁰ [N﹗(C₂OH)][NOSO₃],¹⁰ and [(HOCC₂)C₃im][NTf₂].¹¹ Pal-omar et al. showed the hydroxyl functionalized cations (([HOCC₂)C₃im] or [N₁₁₁₂C₁OH]) with fluorinated anions enhance the solubility compared to nonfunctionalized anologs.⁹ Li et al. also reported that the NH₃ absorption capacity for hydroxyl functionalized ILs ([[(HOCC₂)C₃im][X] X = PF₆, BF₄, DCA, SCN, NO₃] with any given anion was larger compared to the traditional imidazolium-based counterparts ([C₄C₁im][X]).¹² In addition, they also found that ILs with fluorine-containing anions have a higher NH₃ solubility compared with nonfluorinated anions with hydroxyl-containing cations.¹² However, the hydroxyl group on the imidazolium cation led to an increase in viscosity of the ILs.¹⁰,¹¹ Also, contrary to aprotic ILs, the cation chain length for protic ILs was found to have little effect on the NH₃ solubility; however, fluorinated anions such as [NTf₂] had higher NH₃ absorption compared to nonfluorinated anions [SCN] and [NO₃].¹⁶

Thermodynamic models such as the Redlich–Kwong cubic equation of state (EoS),³,⁴ Flory–Huggins model,¹⁷ van der Waals EoS,¹⁸ Peng–Robinson EoS with Kwak and Mansoori mixing rule (PR/KM),¹⁹ Artificial Neural Networks method (AAN),²⁰ modified UNIFAC model,¹³,¹⁴ UNIFAC model,²¹ COSMO-based process simulation with Aspen Plus/Aspen HYSYS,²² Non-Random Two Liquids (NRTL) model,²³ and Antoine equation have been used to correlate experimental VLE data for NH₃ + IL mixtures.

Traditional aprotic ILs (i.e., imidazolium-based ILs) have shown relatively low affinity for NH₃ compared to functionalized ILs, metal containing ILs, and protic ILs. However, the latter suffer from a few major drawbacks. For example, hydroxyl functionalized ILs have higher viscosities, metal containing ILs chemically interact/react with NH₃, and protic ILs are not stable even at room temperature. Therefore, traditional aprotic ILs are still preferable based on reversible NH₃ sorption, lower viscosity, good thermal stability, and lower cost.

The need for accurate thermodynamic and kinetic analysis of NH₃ in ILs is essential for developing existing and new applications. Up to now, thermodynamic measurements for NH₃ in ILs have been conducted using either volumetric or semig gravimetric methods. To the best of our knowledge, the gravimetric method described in this work is one of the most accurate techniques for measuring gas solubility has not been previously utilized to measure the solubility of NH₃ in ILs because NH₃ would permanently damage the balance components. However, the development of the new Hiden XEMIS balance utilizes magnetic coupling to prevent the contact of the gas with the balance electronics so that measurements with NH₃ and other corrosive gases are now possible.²³ In addition to thermodynamic analysis, the XEMIS gravimetric microbalance allows simultaneous time dependent analysis to analyze the kinetics of the dissolution process. Only Bedia et al. have reported the effective diffusivity of NH₃ in one imidazolium and four functionalized ILs using our kinetics model we previously developed for CO₂ + IL mixtures.²⁴ Therefore, in this study, the solubility and diffusivity of NH₃ in [C₄C₁im][PF₆], [C₄C₁im][BF₄], and [C₄C₁im][NTf₂] were analyzed using the gravimetric method at pressures up to 0.7 MPa and at temperatures of 283.15, 298.15, 323.15, and 348.15 K. The accuracy of the experimental method was thoroughly assessed by measuring the solubility of CO₂ in 1-hexyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide ([C₄C₁im][NTf₂]) and compared with the benchmark study developed and sponsored by the International Union of Pure and Applied Chemistry (IUPAC) for the same binary system. Experimental vapor–liquid–equilibrium data were modeled using two approaches: (1) liquid and vapor phase are modeled with the Peng–Robinson EoS, and (2) liquid phase was modeled with an activity coefficient model and the vapor phase fugacity was corrected using the second virial coefficient. In addition to thermodynamics of NH₃ + IL mixtures, the time-dependent behavior of NH₃ absorption in [C₄C₁im][PF₆], [C₄C₁im][BF₄], and [C₄C₁im][NTf₂] were analyzed using the one-dimensional (1D) mass transfer equation and a semi-theoretical Einstein–Stokes equation where the diffusing radius for NH₃ in the mixtures was also reported.

2. EXPERIMENTAL SECTION

2.1. Materials. High purity anhydrous NH₃ (purity of 0.999992%, CAS No. 7664-41-7) and research grade carbon dioxide (purity of 99.9999%, CAS No. 124-38-9) were obtained from Matheson Tri-Gas, Inc. (Topeka, KS). The ILs [C₄C₁im][PF₆] (assay >99%, CAS No. 382150-50-2, Lot No. EQ500831 632), [C₄C₁im][BF₄] (assay ≥96%, CAS No. 174501-64-5, Lot and Filling Code No. 142554 30407904), [C₄C₁im][NTf₂] (assay ≥99%, CAS No. 174899-90-2, IL-0023-HP-1000 H00620.1) were obtained from EMD Chemicals, Inc. (U.S.), Fluka (Switzerland), Fluka (Germany), and IoLiTec, Inc. (Germany), respectively. The chemical structures for the imidazolium-based ILs are provided in Table S1 (see Supporting Information). The ILs were stored under nitrogen to prevent moisture contamination. The as-received ILs were dried under vacuum before conducting gas solubility measurements. To account for buoyancy effects in the experimental method (see Supporting Information), the gas density as a function of T and P as well as the liquid density of the pure ILs as a function of T are required. The density of NH₃ and CO₂ were obtained using the National Institute of Standards and Technology (NIST) REPROP V.9.1 database. The density of [C₄C₁im][NTf₂] was calculated using the recommended IUPAC correlation. The density of [C₄C₁im][PF₆], [C₄C₁im][BF₄], and [C₄C₁im][NTf₂] were obtained from the literature.²⁵

2.2. Apparatus and Measurement. The solubility of NH₃ in the three ILs was measured using a XEMIS gravimetric microbalance (Hiden Isochema, United Kingdom). The experimental apparatus and procedure have been previously discussed in detail;²³ therefore, only brief description along
with measurement uncertainties are provided here. A small amount of IL sample (~40–60 mg) was loaded into a flat bottom Pyrex sample container. The sample was dried and degassed under high vacuum (10^{-12} \text{MPa}) at 348.15 K for 24 h to remove moisture and volatile impurities before each experiment in the balance. The balance was operated in static mode to eliminate drag forces by introducing NH\textsubscript{3} to the top of the balance away from the sample and by controlling the set-point pressure with simultaneous adjustments to the admit and exhaust valves. To ensure sufficient time for thermodynamic equilibrium, the IL sample was maintained at each set point pressure for a minimum of 8 h to a maximum of 20 h. The sample and counterweight temperatures were measured with a K-type thermocouple with an accuracy of ±0.1 K and calibrated against a NIST traceable standard platinum resistance thermometer (Hart Scientific SPRT model 5699 and readout Hart Scientific Blackstack model 1560 with a SPRT module 2560). The Blackstack instrument and SPRT module are also a certified secondary temperature standard with a NIST traceable accuracy to ±0.005 K. The XEMIS microbalance is equipped with two pressure transducers. The low-pressure transducer is used up to 2 MPa, and the high-pressure transducer is used up to 17 MPa. The low-pressure transducer can measure pressures as low as 10 to 20 mbar which is useful when fine pressure regulation is required. Both pressure transducers within the XEMIS microbalance were calibrated against a NIST traceable Paroscientific Model 765IK pressure transducer (range 0 to 6.89 MPa, serial no. 101314). This instrument is a NIST-certified secondary pressure standard with a traceable accuracy of ±0.0008 MPa. The instrumental uncertainty in $T$ is within ±0.1 K, and in $P$ is within ±0.0001 MPa for the low-pressure transducer and ±0.001 MPa for the high-pressure transducer. In this study, NH\textsubscript{3} experiments were conducted using the low-pressure transducer due to the required fine pressure control at low pressures. However, the XEMIS system with both low- and high-pressure transducers was thoroughly evaluated using the CO\textsubscript{2} + [C\textsubscript{6}C\textsubscript{1}im][NTf\textsubscript{2}] reference system described in the IUPAC study.\textsuperscript{26} The gas sorption data were corrected for buoyancy and volume expansion using the procedure described in the Supporting Information (Appendix 1). The total uncertainties in the solubility data have been estimated to be less than ±0.5 mol % at any given $T$ and $P$.

### 2.3. Safety

Ammonia must be handled with extreme caution as it is both a combustible and highly toxic gas. In our laboratory, numerous safety features have been implemented to safely handle NH\textsubscript{3}.\textsuperscript{30} First, the microbalance and temperature/pressure/vacuum control system are located in a specially designed ventilated enclosure which offers protection in the event of an NH\textsubscript{3} leak. The ventilated enclosure is equipped with a fire sprinkler and a safety interlock system. The interlock system is connected to an emergency crash button outside the enclosure, a photoelectric for measuring the ventilation flow rate, and two NH\textsubscript{3} gas monitors (3M Scott Safety, Freedom 5000, Serial 3568 with NH\textsubscript{3} sensor, 096-1965-0100 and 3M Scott Safety, Meridian Universal Gas Detector, Model 096-3480-01 with NH\textsubscript{3} sensor 096-3473-03). In the event that the crash button is depressed, a loss of ventilation occurs, or either gas monitor detects an NH\textsubscript{3} leak, an air-to-open (ATO) valve on the NH\textsubscript{3} feed line will automatically close and shutoff the NH\textsubscript{3} source. One NH\textsubscript{3} monitor is located inside the enclosure and the other NH\textsubscript{3} monitor is located outside the enclosure. Both NH\textsubscript{3} monitors are calibrated and tested quarterly to ensure proper operation. Two alarms were set for alert at 25 ppm (50% of NH\textsubscript{3} allowable exposure limit (AEL)) and warning at 15 ppm (30% NH\textsubscript{3} AEL). In addition, to the ATO valve closing, yellow and red warning lights flash and a siren sounds in the lab for alert and warning alarms, respectively. A data acquisition system (Yokogawa, Model GM10 with Power Supply, Model GM90PS and Module Base, Model GM90MB) provides text and email messages to inform researchers when the NH\textsubscript{3} detector activates or loss of ventilation occurs. Nitrogen gas is setup to purge NH\textsubscript{3} lines and the microbalance (three times) before opening the system. A Research Hazard Review (RHR) was carefully conducted to evaluate and document all hazards including unattended operation, and only authorized operators are allowed to use the equipment. The RHR analysis, procedures, and equipment have demonstrated that the XEMIS microbalance can be safely operated using flammable and toxic gases such as NH\textsubscript{3}.

### 3. RESULTS AND DISCUSSION

#### 3.1. Assessment of Experimental Method

The accuracy of the XEMIS microbalance was assessed by measuring the solubility of CO\textsubscript{2} in [C\textsubscript{6}C\textsubscript{1}im][NTf\textsubscript{2}] which is a standard measurement recommended by the IUPAC in 2007.\textsuperscript{26} The experimental data obtained in this work were compared with the values reported in the IUPAC study by Shiflett et al.\textsuperscript{31} and Kumelan et al.\textsuperscript{32} which are both in good agreement. The solubility of CO\textsubscript{2} in [C\textsubscript{6}C\textsubscript{1}im][NTf\textsubscript{2}] was measured using the high-pressure transducer at pressures from 0.1 to 5 MPa, and using the low-pressure transducer at pressures from 0.01 to 2 MPa. At low pressures, the data showed an average deviation of less than 0.5 mol % compared to Shiflett et al.\textsuperscript{31} At high pressures, the data showed an average deviation of 0.3 mol % compared to Kumelan et al.\textsuperscript{32} The results indicate that the XEMIS gravimetric microbalance is a highly accurate and reliable experimental method for gas solubility measurements in ILs.

The solubility of NH\textsubscript{3} in [C\textsubscript{6}C\textsubscript{1}im][BF\textsubscript{4}] at 293.15 K and at pressures up to 0.5 MPa was selected as a reference system to validate the accuracy of the experimental techniques in the literature. Figure 1 shows this study is in good agreement with

![Figure 1. Comparison of the experimental PTx diagram of NH3 solubility in [C6C1im][BF4] at 293.15 K. Symbols: (○) this study; (△) Li et al.\textsuperscript{33} (◇) Bedia et al.\textsuperscript{34} (V) Palomar et al.\textsuperscript{35} (□) Shang et al.\textsuperscript{36} (☆) Huang et al.\textsuperscript{37} Solid line added to guide the eye.](Image)
Table 1. Henry’s Law Constants for the Mixtures of NH$_3$ + [C$_4$C$_1$im][PF$_6$], NH$_3$ + [C$_4$C$_1$im][BF$_4$], and NH$_3$ + [C$_2$C$_1$im][NTf$_2$] at 283.15, 298.15, 323.15, and 348.15 K

<table>
<thead>
<tr>
<th>Binary System</th>
<th>$T = 283.15$ K</th>
<th>$T = 298.15$ K</th>
<th>$T = 323.15$ K</th>
<th>$T = 348.15$ K</th>
</tr>
</thead>
<tbody>
<tr>
<td>NH$_3$ + [C$_4$C$_1$im][PF$_6$]</td>
<td>0.31 ± 0.02</td>
<td>0.46 ± 0.03</td>
<td>0.83 ± 0.03</td>
<td>1.4 ± 0.16</td>
</tr>
<tr>
<td>NH$_3$ + [C$_4$C$_1$im][BF$_4$]</td>
<td>0.29 ± 0.02</td>
<td>0.41 ± 0.04</td>
<td>0.64 ± 0.07</td>
<td>1.3 ± 0.17</td>
</tr>
<tr>
<td>NH$_3$ + [C$_2$C$_1$im][NTf$_2$]</td>
<td>0.31 ± 0.02</td>
<td>0.46 ± 0.03</td>
<td>0.81 ± 0.04</td>
<td>1.1 ± 0.14</td>
</tr>
</tbody>
</table>

“The uncertainties are the standard error of the coefficient obtained in the linear regression.”

Palomar et al., Bedia et al., and Shang et al. may be due to improper calibration, high relative uncertainty in their given experimental procedure, or the presence of water and/or other impurities in the ILs.

3.2. Vapor Liquid Equilibria Measurement and Thermodynamic Modeling. The solubility of NH$_3$ in [C$_4$C$_1$im][PF$_6$], [C$_4$C$_1$im][BF$_4$], and [C$_2$C$_1$im][NTf$_2$] was measured at temperatures of 283.15, 298.15, 323.15, and 348.15 K and pressures ranging from 0.010 to 0.7 MPa. The present experimental solubility ($P_{Tx}$) data are summarized in Tables S2–S4 (see Supporting Information). Yokozeki and Shiflett previously demonstrated high NH$_3$ sorption in [C$_4$C$_1$im][PF$_6$], [C$_4$C$_1$im][BF$_4$], and [C$_2$C$_1$im][NTf$_2$]; however, the inaccuracy due to weighing small amounts of NH$_3$ using a semigravimetric method resulted in large uncertainties in compositions, especially at low NH$_3$ concentrations (i.e., low pressures). The comparison of the results in this study and our previous results (Figure S1 in the Supporting Information) indicates that the NH$_3$ solubility in these ILs was originally underestimated particularly at low pressures. As described in Section 3.1, the XEMIS balance provides highly accurate solubility measurements even at low pressures (e.g., 10 mbar). Therefore, this study reports the most accurate vapor liquid equilibria data for mixtures of NH$_3$ and [C$_4$C$_1$im][PF$_6$], NH$_3$ and [C$_4$C$_1$im][BF$_4$], and NH$_3$ and [C$_2$C$_1$im][NTf$_2$] at temperatures from 283.15 to 348.15 K and pressures up to 0.7 MPa.

3.2.1. Henry’s Law Constants at Infinite Dilution. Henry’s law constants ($k_H$) are used for assessment of gas solubility in a solvent at dilute concentrations where the lower the $k_H$ value, the higher the gas solubility in the solvent. In this study, the NH$_3$ solubility linearly increases at pressures up to about 0.15 MPa indicating the Henry’s law regime. The results suggest the partial pressure of NH$_3$ ($P_{NH3} \approx P_{t,\text{vap}} \approx 0$) is directly proportional to its liquid phase concentration in the dilute regime. The Henry’s law constants can be obtained from experimental gas solubility ($P_{Tx}$) data using the following relation assuming the hydrostatic pressure correction (Kri-chesky–Karaszynsky equation) is not required:

$$k_H = \lim_{x_i \to 0} \frac{f^V(T, P, y_i)}{x_i} \approx \frac{d f^V}{d x_i} \bigg|_{x_i=0}$$

where $f^V$ is the vapor phase fugacity of the pure gas ($y_1 = 1$) and can be calculated by a proper EOS model at a given condition ($T, P$). The Henry’s law constants were obtained by numerically calculating the limiting slope while $x_{NH3}$ approaches zero using the linear fit of experimental data up to 0.15 MPa including a theoretical zero point (zero pressure and zero composition). The Henry’s law constants obtained for this study are shown in Table 1. The Henry’s law constants increase with an increase in $T$ for the three ILs indicating that the solubility of NH$_3$ decreases with increase in $T$.

The Henry’s law constants are also used to estimate the enthalpy ($\Delta H_{sol}$) and entropy ($\Delta S_{sol}$) of gas dissolution that accompanies the absorption of a mole of gas into a solution at infinite dilution using the thermodynamic van’t Hoff relations:

$$\frac{\partial \ln k_H}{\partial T} _T = -\frac{\Delta H_{sol}}{RT^2}$$

$$\frac{\partial \ln k_H}{\partial T} _T = -\frac{\Delta S_{sol}}{RT}$$

The $\Delta H_{sol}$ and $\Delta S_{sol}$ can be obtained by calculating the slope from eqs 2 and 3 and using the Henry’s law constants provided in Table 1. The $\Delta H_{sol}$ and $\Delta S_{sol}$ of NH$_3$ absorption in [C$_4$C$_1$im][PF$_6$], [C$_4$C$_1$im][BF$_4$], and [C$_2$C$_1$im][NTf$_2$] are reported in Table 2. The negative $\Delta H_{sol}$ indicates the NH$_3$ sorption in imidazolium-based ILs is exothermic. Also, the magnitude of $\Delta H_{sol}$ indicates relatively weak association of NH$_3$ in IL which suggests only physical sorption (i.e., gas solubilities) at low NH$_3$ compositions. The negative $\Delta S_{sol}$ indicates NH$_3$ absorption in imidazolium-based ILs increases in molecular ordering. The $\Delta H_{sol}$ and $\Delta S_{sol}$ results are very similar to SO$_2$ in [C$_4$C$_1$im][BF$_4$] and CO$_2$ in [C$_2$C$_1$im][PF$_6$]. This suggests there is no regular bonding between the solvent (ILs) and solute molecules (NH$_3$, SO$_2$, or CO$_2$) at infinite dilution. It is worth emphasizing that the $\Delta H_{sol}$ and $\Delta S_{sol}$ results are only instructive at low NH$_3$ compositions (i.e., below 0.15 MPa).

3.2.2. Peng–Robinson Equation of State Model. The Peng–Robinson Equation of State (PR-EoS) was selected to model vapor–liquid equilibrium data, as it has been used extensively in academic evaluations and industrial applications and has been previously applied to NH$_3$ + IL mixtures. The PR-EoS is given as:

$$P = \frac{RT}{V - b_m} - \frac{a_m}{V(V - b_m) + b(V - b_m)}$$
where $a_m$ and $b_m$ are the mixture attractive term and co-volume parameters, respectively. The $a_m$ and $b_m$ parameters are computed with the one parameter van der Waals mixing rule with Boston–Mathias extension for the attractive term and the standard PR-EoS alpha function. Pure component parameters ($a_i$ and $b_i$) are computed based on the critical temperature ($T_c$), critical pressure ($P_c$), and acentric factor ($\omega$). The PR-EoS parameters ($a_m$, $b_m$, $\alpha(T)$, $m$, $n_{av}$, $k_{ij}$, and $l_{ij}$) used for this analysis are shown in eqs 5–12. Table 3 provides the EoS constants used for NH$_3$ and ILs. The binary parameters $k_{ij}$ and $l_{ij}$ were determined from the regression of the VLE data using ASPEN Plus v.1036 by minimizing the objective function of the average absolute relative deviation (% AARD = (100/n)\sum_{i=1}^{n} |(P_{exp} - P_{regressed})/P_{exp})|). The results are listed in Table 4. The experimental results along with the PR-EoS models are shown in Figures 2, 3, and 4. The difference between the experimental and estimated VLE data for the three binary mixtures of NH$_3$–ILs was less than 5% AARD except in the case of NH$_3$ + [C$_4$C$_1$im][PF$_6$] at 298.15 K where the model underestimates the VLE behavior by 6.5% AARD. The critical properties of the ILs cannot be experimentally determined (i.e., critical temperature is above the decomposition temperature). Therefore, the critical properties used in this study were obtained from the literature where the critical properties were estimated using a Group Contribution Method. The critical properties and acentric factor of the ILs were taken from ref 37.

Table 3. Pure Component EoS Constants Used in This Study

<table>
<thead>
<tr>
<th>Compound</th>
<th>Molar Mass (g mol$^{-1}$)</th>
<th>$T_c$ (K)</th>
<th>$P_c$ (MPa)</th>
<th>$\omega$</th>
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<tr>
<td>NH$_3$</td>
<td>17.03</td>
<td>405.7</td>
<td>11.3</td>
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<tr>
<td>[C$_4$C$_1$im][PF$_6$]</td>
<td>284.18</td>
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<td>[C$_4$C$_1$im][BF$_4$]</td>
<td>226.03</td>
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<td>[C$_4$C$_1$im][NTf$_2$]</td>
<td>391.31</td>
<td>1244.9</td>
<td>3.26</td>
<td>0.182</td>
</tr>
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</table>

Table 4. Binary Interaction Parameters for the Peng–Robinson EoS Model

<table>
<thead>
<tr>
<th>Compound</th>
<th>Temp (K)</th>
<th>$k_{12}$</th>
<th>$l_{12}$</th>
<th>AARD %</th>
</tr>
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<tbody>
<tr>
<td>NH$_3$-[C$_4$C$_1$im][PF$_6$]</td>
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<tr>
<td>298.15</td>
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<tr>
<td>323.15</td>
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<td>348.15</td>
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<td>NH$_3$-[C$_4$C$_1$im][BF$_4$]</td>
<td>283.15</td>
<td>$-0.2034$</td>
<td>$-0.1648$</td>
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</tr>
<tr>
<td>298.15</td>
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<tr>
<td>323.15</td>
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<tr>
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</tbody>
</table>

Figure 2. PT$x$ phase diagram for NH$_3$ and [C$_4$C$_1$im][PF$_6$] at 283.15, 298.15, 318.15, and 338.15 K. Symbols represent experimental data, and solid lines represents PR-EoS model.

Figure 3. PT$x$ phase diagram for NH$_3$ and [C$_4$C$_1$im][BF$_4$] at 283.15, 298.15, 318.15, and 338.15 K. Symbols represent experimental data, and solid lines represents PR-EoS model.

Figure 4. PT$x$ phase diagram for NH$_3$ and [C$_4$C$_1$im][NTf$_2$] at 283.15, 298.15, 318.15, and 338.15 K. Symbols represent experimental data, and solid lines represents PR-EoS model.
respectively; therefore, the results indicate that the $T_i$ calculated in Table 3 can be used to correlate NH$_3$ + IL systems using the PR-EoS model.

**Pure component parameters:**

$$a_i = \frac{0.45724RT_i^2}{P_i} \alpha_i(T)$$

(5)

$$b_i = \frac{0.07780RT_i}{P_i}$$

(6)

**Alpha Function:**

$$\alpha_i(T) = (1 + m_i(1 - \sqrt{T_i}))^2$$

(7)

$$m_i = 0.37464 + 1.54226\omega_i - 0.26992\omega_i^2$$

(8)

$$T_i = \frac{T}{T_i}$$

(9)

**Mixing rule:**

$$a_m = \sum_i \sum_j x_j (1 - k_{ij})(a_i a_j)^{1/2}$$

$$+ \sum_i \left( \sum_j x_j (a_i a_j)^{1/2} l_{ij} \right)^3$$

(10)

$$k_{ij} = k_{ij}^{(1)} + k_{ij}^{(2)} T + \frac{k_{ij}^{(3)}}{T}$$

and

$$l_{ij} = l_{ij}^{(1)} + l_{ij}^{(2)} T + \frac{l_{ij}^{(3)}}{T}$$

(11)

$$b_m = \sum_i x_i b_i$$

(12)

**Activity Coefficient Model.** At low to medium pressures, EoS models can be used to estimate the vapor phase fugacity as nonidealities in the vapor phase are usually small. On the other hand, nonidealities in the liquid phase can be large; therefore, the EoS with van der Waals one fluid mixing rules may not reasonably predict the liquid phase fugacities.38 Alternatively, EoS models with excess free energy based mixing rules can be used to estimate the liquid phase fugacities at all temperatures and pressures.38 However, modeling the experimental data with EoS models is usually more tedious while activity coefficient models are simple to estimate the liquid phase fugacities at low pressures. When two different models are used (an activity coefficient model for the liquid phase and an EoS model for the vapor phase), the critical point of the mixture might be incorrectly predicted because the properties of the two phases might not be identical.38 However, considering the critical point of the mixture is considerably high, using two different models can accurately estimate the VLE of NH$_3$ and IL mixtures. Therefore, in this study, the vapor liquid equilibria of NH$_3$ + IL mixtures were also modeled using the Non-Random Two Liquid (NRTL) solution model for the liquid phase and the Second Virial Coefficient correction for the vapor phase.

For low and medium pressure, VLE for an N-component system can be described:24

$$y_i P_f = x_i P_i^\text{nop}$$

(13)

where $y_i$ is the vapor phase mole fraction for $i$th species, $x_i$ is the liquid phase mole fraction for $i$th species, $P$ is pressure, $P_i^{\text{nop}}$ is the saturated vapor pressure for $i$th species, $\phi_i$ is a Poynting pressure correction factor for $i$th species of a gas, and $\gamma_i$ is the activity coefficient for $i$th species (function of composition at $T$). For a binary system of NH$_3$ + IL mixtures, it is reasonable to assume that the solubility of IL in NH$_3$ is negligible ($P_i^{\text{nop}} \approx 0$) so $y_{NH_3} = 1$ (or $y_{IL} = 0$). The activity coefficient of NH$_3$ (1) is given by

$$\gamma_i = \frac{P_f x_i}{P_i^{\text{nop}}}$$

(14)

The Poynting pressure correction factor ($\phi_i$) for NH$_3$ in the present case is

$$\phi_i = \exp \left( \frac{(B_1 - \bar{V}_i)(P - P_i^{\text{nop}})}{RT} \right)$$

(15)

where $B_1(T)$ is the second virial coefficient of NH$_3$ at system $T$, $\bar{V}_i$ is the saturated molar liquid volume at system $T$, and $R$ is the universal gas constant. The $B_1(T)$ is obtained using the NIST REFPROP computer code v.9.1.25 $\bar{V}_i$ can be calculated as described in the Supporting Information, and $P_i^{\text{nop}}$ can be obtained using the Extended Antoine vapor pressure model:

$$\ln P_i^{\text{nop}} = A + \frac{B}{T + C} + DT + E \ln T + FT^G$$

(16)

where $T(K), P_i^{\text{nop}} (kPa)$, and coefficients are given as $A = 83.58, B = -4669.70, C = 0, D = 0, E = -11.61, F = 0.02$, and $G = 1.36$ The activity coefficients were calculated using the NRTL equations:

$$\ln \gamma_i = x_2^2 \left[ \tau_{12} \left( \frac{G_{21}}{x_1 + x_2 G_{21}} \right)^2 + \frac{\tau_{12} G_{12}}{(x_1 + x_2 G_{12})^2} \right]$$

(17)

$$\ln \gamma_2 = x_1^2 \left[ \tau_{12} \left( \frac{G_{21}}{x_1 + x_2 G_{21}} \right)^2 + \frac{\tau_{12} G_{12}}{(x_1 + x_2 G_{12})^2} \right]$$

(18)

where $G_{12}$ and $G_{21}$ are defined by the interaction parameters ($\tau_{12}, \tau_{12}$, $\alpha$):

$$G_{12} \equiv \exp(-\alpha r_{12})$$

(19)

and

$$G_{21} \equiv \exp(-\alpha r_{21})$$

where the nonrandomness factor ($\alpha$) is assumed to be a constant of 0.2. In this study, binary interaction parameters were obtained using only temperature-dependent terms:

$$r_{12} = r_{12}^{(0)} + \frac{r_{12}^{(1)}}{T(K)}$$

(20)

and

$$r_{21} = r_{21}^{(0)} + \frac{r_{21}^{(1)}}{T(K)}$$

(21)

The regressed binary interaction parameters are summarized in Table 5.

The experimental VLE results and NRTL models are shown in Figures 5, 6, and 7. The models accurately predict the VLE for each of the three systems with an average overall difference of less than 5% AARD. The activity coefficient model for the

### Table 5. Binary Interaction Parameters for the NRTL Activity Coefficient Model

<table>
<thead>
<tr>
<th>System (1)/(2)</th>
<th>$r_{12}^{(0)}$</th>
<th>$r_{12}^{(1)}$</th>
<th>$r_{21}^{(0)}$</th>
<th>$r_{21}^{(1)}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>NH$_3$ + [C$_4$C$_1$im][NTf$_2$]</td>
<td>-3.306</td>
<td>1669.7</td>
<td>-0.1201</td>
<td>-730.79</td>
</tr>
<tr>
<td>NH$_3$ + [C$_4$C$_1$im][PF$_6$]</td>
<td>-2.439</td>
<td>1350.8</td>
<td>0.3428</td>
<td>-815.96</td>
</tr>
<tr>
<td>NH$_3$ + [C$_4$C$_1$im][BF$_4$]</td>
<td>-4.871</td>
<td>2634.5</td>
<td>-0.1423</td>
<td>-891.45</td>
</tr>
</tbody>
</table>
NH₃ and [C₄C₁im][PF₆] system is also compared with results reported by Tomida et al. The activity coefficient model accurately estimated the NH₃ composition at high pressures using only low pressure (up to 0.7 MPa) VLE measurements. The average absolute relative deviation between this study and Tomida et al. was 5.1%, 3.5%, 2.5%, and 2.8% AARD at 283.15, 298.15, 323.15, and 348.15 K, respectively. The excellent agreement between the model obtained in this study and the high pressure experimental data obtained by Tomida et al. indicates the low pressure solubility measurements were highly accurate. The models also do not predict any liquid–liquid phase separation at high NH₃ concentration.

Chemical interaction between NH₃ and imidazolium-based ILs was also considered by measuring the desorption of [C₄C₁im][PF₆], [C₄C₁im][BF₄], and [C₂C₁im][NTf₂] at temperatures of 283.15, 298.15, 323.15, and 348.15 K and at pressures ranging from 0.010 to 0.7 MPa. The desorption data are shown in Figures 5, 6, and 7. At some conditions, particularly at low T and low P, the difference between absorption and desorption solubility was as high as ~3 mol %. This difference between the equilibrium concentrations for the given isotherms during absorption and desorption suggests the interaction between NH₃ and imidazolium-based ILs may be both chemical and physical. The interaction of NH₃ and ILs requires more in-depth spectroscopic analysis which is currently under investigation.

3.2.4. Comparison of NH₃ + [C₄C₁im][PF₆], NH₃ + [C₄C₁im][BF₄], and NH₃ + [C₄C₁im][NTf₂] Systems. The solubility of NH₃ in [C₄C₁im][PF₆], [C₄C₁im][BF₄], and [C₂C₁im][NTf₂] were also compared at 298.15 K. As can be seen in Figure 8, all three ILs have very similar NH₃ sorption capacities even though there is a slight difference between the solubilities up to ~0.72 mole fraction NH₃ following the order: [C₄C₁im][PF₆] < [C₂C₁im][NTf₂] < [C₄C₁im][BF₄]. Above 0.72 mole fraction NH₃ all three ILs demonstrated similar solubility behavior at 298.15 K. The [C₄C₁im][PF₆] and [C₂C₁im][BF₄] ILs share the same cation and show a slight difference in solubility; therefore, the anion may have only a minor impact on the NH₃ solubility. If we assume the anion has a minor effect on solubility, then the shorter alkyl chain

NH₃ and [C₄C₁im][PF₆] system is also compared with results reported by Tomida et al. The activity coefficient model accurately estimated the NH₃ composition at high pressures using only low pressure (up to 0.7 MPa) VLE measurements. The average absolute relative deviation between this study and Tomida et al. was 5.1%, 3.5%, 2.5%, and 2.8% AARD at 283.15, 298.15, 323.15, and 348.15 K, respectively. The excellent agreement between the model obtained in this study and the high pressure experimental data obtained by Tomida et al. indicates the low pressure solubility measurements were highly accurate. The models also do not predict any liquid–liquid phase separation at high NH₃ concentration.

Chemical interaction between NH₃ and imidazolium-based ILs was also considered by measuring the desorption of [C₄C₁im][PF₆], [C₄C₁im][BF₄], and [C₂C₁im][NTf₂] at temperatures of 283.15, 298.15, 323.15, and 348.15 K and at pressures ranging from 0.010 to 0.7 MPa. The desorption data are shown in Figures 5, 6, and 7. At some conditions, particularly at low T and low P, the difference between absorption and desorption solubility was as high as ~3 mol %. This difference between the equilibrium concentrations for the given isotherms during absorption and desorption suggests the interaction between NH₃ and imidazolium-based ILs may be both chemical and physical. The interaction of NH₃ and ILs requires more in-depth spectroscopic analysis which is currently under investigation.

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Table 6. Average Effective Diffusion Coefficients for NH$_3$ in [C$_4$C$_1$im][PF$_6$], [C$_4$C$_1$im][BF$_4$], and [C$_2$C$_1$im][NTf$_2$] Systems

<table>
<thead>
<tr>
<th>Ionic Liquid</th>
<th>Diffusivity ($\times 10^{-10}$ m$^2$·s$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>[C$_4$C$_1$im][PF$_6$]</td>
<td>1.9 ± 0.13</td>
</tr>
<tr>
<td>[C$_4$C$_1$im][BF$_4$]</td>
<td>1.8 ± 0.09</td>
</tr>
<tr>
<td>[C$_2$C$_1$im][NTf$_2$]</td>
<td>2.8 ± 0.25</td>
</tr>
<tr>
<td>[C$_4$C$_1$im][PF$_6$]</td>
<td>3.1 ± 0.07</td>
</tr>
<tr>
<td>[C$_4$C$_1$im][BF$_4$]</td>
<td>3.3 ± 0.13</td>
</tr>
<tr>
<td>[C$_2$C$_1$im][NTf$_2$]</td>
<td>5.2 ± 0.16</td>
</tr>
<tr>
<td>[C$_4$C$_1$im][PF$_6$]</td>
<td>4.0 ± 0.04</td>
</tr>
<tr>
<td>[C$_4$C$_1$im][BF$_4$]</td>
<td>4.8 ± 0.05</td>
</tr>
<tr>
<td>[C$_2$C$_1$im][NTf$_2$]</td>
<td>8.5 ± 0.65</td>
</tr>
<tr>
<td>[C$_4$C$_1$im][PF$_6$]</td>
<td>7.1 ± 0.22</td>
</tr>
<tr>
<td>[C$_4$C$_1$im][BF$_4$]</td>
<td>6.0 ± 0.26</td>
</tr>
<tr>
<td>[C$_2$C$_1$im][NTf$_2$]</td>
<td>18.3 ± 2.1</td>
</tr>
</tbody>
</table>

*The uncertainties are due to the random errors as a result of mass measurement in the balance and systematic error as a result of change in L.

length may lead to an increase in NH$_3$ solubility ([C$_4$C$_1$im][PF$_6$] < [C$_4$C$_1$im][NTf$_2$]). Figure 8 also shows the large negative deviation from ideal solubility behavior (Raoult’s law) indicating a strong NH$_3$–IL interaction.

### 3.3. Absorption Kinetics and Diffusion Coefficients

In addition to the thermodynamic analysis of NH$_3$ and IL mixtures, the time-dependent behavior of NH$_3$ dissolution in ILs was analyzed with a simplified Fickian diffusion model. Detailed procedures of the diffusivity analysis are discussed by Shiflett and Yokozeki. The simplification model, the following assumptions are made: (i) gas dissolves through a one-dimensional (vertical) diffusion process, and there is no convective flow in the liquid; (ii) the interaction between gas and IL is physical; (iii) a thin boundary layer between the IL and gas phase exists, and the layer reaches a saturation concentration ($C_s$) at any given temperature and pressure (Boundary Condition 1); (iv) temperature and pressure are constant (experimental design); (v) the NH$_3$–IL solution is dilute, and the thermophysical properties are constant at a given $T$ and $P$ condition; and (vi) gas does not penetrate through the Pyrex cup (Boundary Condition 2). These assumptions lead to the dissolution of gas in the IL for one-dimensional (1D) mass diffusion due to the local concentration difference:

\[
\frac{\partial^2 C}{\partial z^2} = \frac{\partial C}{\partial t} \quad (21)
\]

where $C$ is the concentration of NH$_3$ in IL, $t$ is the time, $z$ is the vertical location, and $L$ is the depth of IL in the sample container. In this study, the depth ($L$) is estimated from the solution mass, the dimension of the sample cup, and the weight fraction averaged density of the solution at initial and final composition for a given $T$ and $P$. Equation 21 can be analytically solved using a separation of variables technique and applying the proper initial and boundary conditions (eqs 22–24), to obtain the concentration profile in the $z$ direction (eq 25).

Initial Condition: \( t = 0 \) \quad 0 < z < L \quad C = C_o \quad (22)\n
Boundary Condition 1: \( t > 0 \) \quad z = 0 \quad C = C_s \quad (23)\n
Boundary Condition 2: \( t > 0 \) \quad z = L \quad \frac{\partial C}{\partial z} = 0 \quad (24)\n
\[ C = C_o \left[ 1 - \left(1 - \frac{C_o}{C_s}\right) \sum_{n=0}^{\infty} \exp(-\lambda_n^2Dt) \sin \lambda_n z \right] L^2 \lambda_n^2 \quad (25)\n
The concentration obtained in the microbalance is the average concentration at a given time, not the concentration profile in $z$; therefore, the space averaged concentration at a given time can be obtained using the following equations:

\[ \langle C \rangle = \int_0^L \frac{C}{L} \, dz \quad (26)\n
\[ \langle C \rangle = C - 2 \left(1 - \frac{C_o}{C_s}\right) \sum_{n=0}^{\infty} \exp(-\lambda_n^2Dt) \sin n \pi \frac{L}{L}\n
\]

Although eq 27 has an infinite summation term, only the first few terms are required for most analysis and the summation term was terminated when the numerical contribution is infinitely small. At any given $T$ and $P$, the rate of gas dissolution at each time step was obtained using eq 27 to obtain $D$ and $C_s$. The effective $D$ value at each $P$ for a given isotherm was obtained by averaging the calculated $D$ obtained using a constant height of the solution at initial ($C_0$) and final ($C_s$) saturation compositions. The analysis of eq 27 was performed using nonlinear regression in MATLAB. Tables S2–S4 in the Supporting Information summarizes the results for $D$ and $C_s$ for NH$_3$ + [C$_4$C$_1$im][PF$_6$], NH$_3 + [C_4C_1im][BF_4]$, and NH$_3 + [C_2C_1im][NTf_2]$ in the effective $D$ for each system as a function of temperature are summarized in Table 6. As expected, the diffusivity of NH$_3$ in [C$_4$C$_1$im][PF$_6$], [C$_4$C$_1$im][BF$_4$], and [C$_2$C$_1$im][NTf$_2$] increases with an increase in temperature as the viscosity of the solution decreases. However, at constant $T$, the pressure dependence of $D$ is rather weak compared to temperature. A key finding in this part of the study is that the diffusivity of NH$_3$ in [C$_4$C$_1$im][PF$_6$], [C$_4$C$_1$im][BF$_4$], and [C$_2$C$_1$im][NTf$_2$] is about 3 to 5 times lower than the effective $D$ in water.

For example, the diffusivity of NH$_3$ in water at 298.15 K is 16 $\times 10^{-10}$ m$^2$·s$^{-1}$ whereas the diffusivity of NH$_3$ in [C$_4$C$_1$im][PF$_6$], [C$_4$C$_1$im][BF$_4$], and [C$_2$C$_1$im][NTf$_2$] at 298.15 K is 3.1 $\times 10^{-10}$, 3.3 $\times 10^{-10}$, and 5.2 $\times 10^{-10}$, respectively.

One way to test the model validity is to check the NH$_3$ solubility ($C_s$) determined from the present analysis. If the model is physically meaningful, the $C_s$ values should be consistent with the experimental solubility values. In fact, the difference between the experimental solubility and model $C_s$ values was less than ~1 mol % for all systems which indicates the model reasonably predicts the $C_s$ and $D$ values. A few comments also should be made regarding the $D$ values. Contrary to the main assumptions of the model, the NH$_3$ + IL mixtures cannot be considered as dilute solution, and consequently the diffusion coefficients can depend on the concentration. In reality, the thermophysical properties of the mixture change upon gas dissolution, and $L$ varies with the amount of gas dissolved in the IL. Therefore, when we apply the present model, the analyzed diffusion coefficients must be regarded as “effective” or “apparent” diffusion coefficients.

#### 3.3.1. Stokes–Einstein Model

When a solute sphere (with radius $r_s$) moves through a continuum fluid, the Stokes–Einstein equation correlates the diffusion coefficient ($D_{AB}$) and the viscosity of solvent ($\mu_B$) assuming the diffusing particle is perfectly spherical where $k$ is the Boltzmann constant and $T$ is the temperature:

\[ \langle C \rangle = C - 2 \left(1 - \frac{C_o}{C_s}\right) \sum_{n=0}^{\infty} \exp(-\lambda_n^2Dt) \sin n \pi \frac{L}{L}\n
\]

\[ \langle C \rangle = C - 2 \left(1 - \frac{C_o}{C_s}\right) \sum_{n=0}^{\infty} \exp(-\lambda_n^2Dt) \sin n \pi \frac{L}{L}\n
\]

\[ \langle C \rangle = C - 2 \left(1 - \frac{C_o}{C_s}\right) \sum_{n=0}^{\infty} \exp(-\lambda_n^2Dt) \sin n \pi \frac{L}{L}\n
\]
The coeﬃcients for Equation 33

<table>
<thead>
<tr>
<th>Compound</th>
<th>(A_i)</th>
<th>(B_i) (K)</th>
<th>(C_i) (K(^{-1}))</th>
<th>(D_i) (K(^{-2}))</th>
</tr>
</thead>
<tbody>
<tr>
<td>NH(_3)</td>
<td>(-5.130771)</td>
<td>922.2</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>([C_2C_1im][PF_6])</td>
<td>(-182.774)</td>
<td>24992.4</td>
<td>(4.84019 \times 10^{-1})</td>
<td>(-4.44779 \times 10^{-4})</td>
</tr>
<tr>
<td>([C_2C_1im][BF_4])</td>
<td>(-149.99)</td>
<td>20757.8</td>
<td>(3.91576 \times 10^{-1})</td>
<td>(-3.55363 \times 10^{-4})</td>
</tr>
<tr>
<td>([C_2C_1im][NTf_2])</td>
<td>(-60.707)</td>
<td>9364.9</td>
<td>(1.49780 \times 10^{-1})</td>
<td>(-1.33200 \times 10^{-4})</td>
</tr>
</tbody>
</table>

“Viscosity in mPa·s (or cP), and T in K.” Parameters obtained using linear ﬁtting of viscosity data from ref 25. Parameters are taken from ref 41.

Table 8. Determined Parameters for Equations 30 and 32

<table>
<thead>
<tr>
<th>System</th>
<th>(a) (ln m(^2) s(^{-1}) K(^{-1}))</th>
<th>(b) (ln m(^2) s(^{-1}) K(^{-1}))</th>
<th>(c)</th>
<th>Radius (nm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>NH(_3) + ([C_4C_1im][PF_6])</td>
<td>(-25.880 \pm 0.186)</td>
<td>(0.466 \pm 0.046)</td>
<td>(0.502 \pm 0.029)</td>
<td>(0.127 \pm 0.024)</td>
</tr>
<tr>
<td>NH(_3) + ([C_4C_1im][BF_4])</td>
<td>(-26.142 \pm 0.253)</td>
<td>(0.504 \pm 0.095)</td>
<td>(0.438 \pm 0.034)</td>
<td>(0.165 \pm 0.042)</td>
</tr>
<tr>
<td>NH(_3) + ([C_4C_1im][NTf_2])</td>
<td>(-25.742 \pm 0.240)</td>
<td>(0.458 \pm 0.093)</td>
<td>(0.899 \pm 0.021)</td>
<td>(0.111 \pm 0.027)</td>
</tr>
</tbody>
</table>

Errors are the standard uncertainty obtained in regression analysis.

The mixture viscosity for an N-component solution can be estimated using the following model:

\[
D = \frac{kT}{6\pi r^3\mu}
\]

The linearized form of eq 29 can be written as

\[
\ln(D/T) = a - b \ln(\mu/\mu_0)
\]

where \(D\) is diffusivity (m\(^2\) s\(^{-1}\)), \(k\) is the Boltzmann constant, \(T\) is temperature (K), and \(\mu_0\) is a unit viscosity (1 mPa·s) that is used as a normalization factor to have proper dimension in the equation. \(a = \ln(k/6\pi r^3\mu_0)\) and \(b\) are the adjustable parameters.

Figure 9. Diffusivity of NH\(_3\) in \([C_4C_1im][PF_6]\). Lines represent model calculations, and symbols represent experimental data: (●) 283.15 K; (◇) 298.15 K; (◆) 323.15 K; (▲) 348.15 K.

4. CONCLUSION

The first vapor liquid equilibrium (VLE) measurements for the binary systems of NH\(_3\) and three imidazolium-based ILs have been successfully measured using a gravimetric microbalance. The vapor liquid equilibrium (VLE) for the binary systems of NH\(_3\) and imidazolium-based ILs \([C_4C_1im][PF_6]\), \([C_4C_1im][BF_4]\), and \([C_4C_1im][NTf_2]\) were measured at 283.15, 298.15, 323.15, and 348.15 K with pressures up to 7 MPa using a...
XEMIS gravimetric microbalance. The thermodynamic modeling using the Peng–Robinson equation of state as well as the Non-Random Two Liquid (NRTL) activity coefficient model was in excellent agreement with the experimental data. The Fickian diffusion coefficients for NH₃ in imidazolium-based ILs were obtained using a one-dimensional mass diffusion equation and found to be about 3 to 5 times lower than the diffusion of NH₃ in water. A semitheoretical Stokes–Einstein equation was used to model diffusivities and to obtain the diffusion radius of NH₃ in imidazolium-based ILs. The Stokes–Einstein equation is useful for correlating the diffusivity of NH₃ in imidazolium-based ILs when the thermophysical properties of the solution are known or properly estimated.

ASSOCIATED CONTENT

Supporting Information
The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.iecr.9b00274.

Experimental data reduction; comparison of PTx diagram NH₃ and [C₄C₁im][PF₆], [C₄C₁im][BF₄], and [C₂C₁im][NTf₂] (PDF)

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Notes
The authors declare no competing financial interest.

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ABBREVIATIONS

CO₂ = carbon dioxide
H₂O = water
NH₃ = ammonia
[C₄C₁im][BF₄] = 1-ethyl-3-methylimidazolium tetrafluoroborate
[C₄C₁im][PF₆] = 1-ethyl-3-methylimidazolium hexafluorophosphate
[C₂C₁im][NTf₂] = 1-ethyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide
[C₂C₁im][NO₃] = 1-ethyl-3-methylimidazolium nitrate
[C₄C₁im][BF₄] = 1-butyl-3-methylimidazolium tetrafluoroborate
[C₄C₁im][PF₆] = 1-butyl-3-methylimidazolium hexafluorophosphate
[C₄C₁im][NTf₂] = 1-butyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide
[C₄C₁im][BF₄] = 1-butyl-3-methylimidazolium tetrafluoroborate
[C₆C₁im][BF₄] = 1-hexyl-3-methylimidazolium tetrafluoroborate
[NH₄]₂[C₂OH][NTf₂] = choline bis(trifluoromethylsulfonyl)imide
[(HOC₂)C₁im][BF₄] = 1-(2-hydroxyethyl)-3-methylimidazolium tetrafluoroborate
[(HOC₂)C₁im][DCA] = 1-(2-hydroxyethyl)-3-methylimidazolium dicyanamide
[(HOC₂)C₁im][NTf₂] = 1-(2-hydroxyethyl)-3-methylimidazolium bis(trifluoromethylsulfonyl)imide
[(HOC₂)C₁im][NO₃] = 1-(2-hydroxyethyl)-3-methylimidazolium nitrate

Figure 10. Diffusivity of NH₃ in [C₄C₁im][BF₄]. Diffusivity vs composition. Lines represent model calculations, and symbols represent experimental data: (■) 283.15 K; (●) 298.15 K; (◆) 323.15 K; (▲) 348.15 K.

Figure 11. Diffusivity of NH₃ in [C₂C₁im][NTf₂]. Lines represent model calculations, and symbols represent experimental data: (■) 283.15 K; (●) 298.15 K; (◆) 323.15 K; (▲) 348.15 K.
Solubilities of Ammonia in Basic Imidazolium Ionic Liquids.

Phase Equilib. (9), 2414–2421.


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